

Polyhalogen Anions

DOI: 10.1002/anie.201209928

Structural Evidence for Undecabromide [Br₁₁]^{-**}

Heike Haller, Jan Schröder, and Sebastian Riedel*

Dedicated to Professor Werner Uhl on the occasion of his 60th birthday

Among all known polyhalides, the chemistry of the iodine compounds has long been known and has been investigated extensively. Already in 1870 Jørgensen^[1] discovered that triiodide was formed by mixing potassium iodide with iodine in an alcoholic solution. To date many examples of different mono-, di-, tri-, and tetraanions are known. [2] Focusing on the polyiodide monoanions, only compounds up to $[X_9]^-$ have to the best of our knowledge been structurally characterized. In 2011 Groessl et al. observed the first higher polyiodides $[I_{11}]^-$, $[I_{13}]^-$, and $[I_{15}]^-$ by mass spectrometry. [3] They mixed [PMIM]I (1-propyl-3-methylimidazolium) with the appropriate amount of iodine, diluted it in acetonitrile, and analyzed this mixture by ESI-MS. To confirm the observed fragmentation pathways, they calculated the stable structures for $[I_{2n+1}]^-$ (n=1-7) at the DFT level. The observed compounds in this gas-phase experiment have the highest ratio of elements to charge of any previously studied polyhalides.

Regarding the lighter and more reactive homologue of the halogens, bromine, many fewer examples are known. Linear $[Br_4]^{2-,[4]}$ Z-shaped $[Br_8]^{2-,[5]}$ ring-shaped $[Br_{10}]^{2-,[6]}$ and the recently found $[Br_{20}]^{2-[7,8]}$ are examples of known polybromide dianions.^[9] Among the polybromide monoanions, only [Br₃] has been known for a long time. Indeed, Chattaway and Hoyle^[10] discussed already in 1923 the existence of higher polybromide monoanions but until 2010 there has been neither vibrational spectroscopy nor X-ray diffraction evidence for higher polybromides, such as [Br₅]⁻, [Br₇]⁻, and [Br₉]^{-.[11]} In 2011 our group reported the evidence for the structure of a higher polybromide monoanion [NPr₄][Br₉].^[12] Shortly thereafter Feldmann et al. synthesized and characterized [P(Ph)₃Br][Br₇] in ionic liquids, [8] and very recently Himmel et al. reported on the structure of [Br₅]⁻.^[13] Beyond these polybromide anions also polychloride and fluoride anions have been characterized by IR and Raman spectroscopy. In the case of polychloride monoanions, [Cl₃]-,[14] $[Cl_3\cdots Cl_2]^{-}$, $[Cl_5]^{-}$, [I6] and $[Cl_9]^{-[16]}$ are known, whereby the latter exists only under cryogenic conditions. This holds also for the only example of a polyfluoride anion $[F_3]^-$, which

[*] Dipl.-Chem. H. Haller, B. Sc. J. Schröder, Dr. S. Riedel Institut für Anorganische und Analytische Chemie Albert-Ludwigs Universität Freiburg Albertstrasse 21, Freiburg (Germany) E-mail: sebastian.riedel@psichem.de Homepage: http://www.psichem.de

[**] This work has been supported by the Fonds der Chemischen Industrie (FCI). We are grateful to Prof. Ingo Krossing and Prof. Harald Hillebrecht for their generous and continuous support.

Supporting information for this article is available on the WWW under http://dx.doi.org/10.1002/anie.201209928.

was characterized by matrix-isolation spectroscopy in argon and neon matrices at 4 K.[17]

Herein, we report the largest structurally known polyhalide monoanion undecabromide $[Br_{11}]^-$, which was recently predicted to be thermochemically stable by 297.0 kJ mol⁻¹ against the elimination of five Br₂ molecules at SCS-MP2/ def2-TZVPP level. [12] The reaction of bis(triphenylphosphine)iminium bromide (PPNBr) with an excess of bromine led to [PPN][Br₁₁·Br₂], which is the highest known ratio of elements to charge of any structural known polyhalide compound.

The single-crystal X-ray structure shows that the salt [PPN][Br₁₁·Br₂] crystallizes in the triclinic space group $P\bar{1}$. Similar to other polyhalides, [Br₁₁·Br₂] is built of the two building blocks Br⁻ and Br₂. A central Br⁻ is coordinating five Br₂ molecules (Figure 1). The charge of the central Br⁻ is

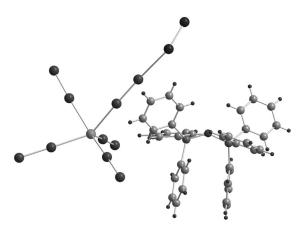


Figure 1. Molecular structure of [PPN][Br₁₁·Br₂] in the solid state. Five bromine molecules are coordinated to a central Br to form the polybromide anion [Br₁₁]⁻.

donated into the antibonding LUMOs of the end-on coordinated Br₂ molecules, thereby weakening and elongating their Br-Br bond by an average of 5.9 pm compared to molecular Br₂.[18]

The $[Br_{11}]^-$ units are linked to form chains in which a sixth bromine molecule is embedded. These chain-linking Br-Br contacts are much longer than the contacts in the $[Br_{11}]^-$ units but nevertheless shorter than twice the bromine van der Waals radius of 370 pm. The contacts between the $[Br_{11}]^{-}$ units and the embedded bromine molecule is also shorter than 370 pm, but the embedded Br₂ molecule is not linked to a Br⁻ anion and the Br-Br distance (0.31 pm) is practically not elongated (Table 1, Figure 2). Therefore this Br₂ molecule is not attributed to the polyhalide monoanion.



Table 1: Measured and computed distances [pm] of $[Br_{11} \cdot Br_2]^{-}$. [a]

Distance	Exp.	SCS-MP2 ^[b]	BP86	PBE0 ^[b]	BHLYP ^[b]
Br1-Br2	294.55(3)	296.2/301.4	292.0	293.1/295.0	303.6/304.7
Br2-Br3	233.31(3)	234.0/238.6	241.7	234.5/236.4	233.6/235.7
Br1-Br4	308.09(3)	299.1/305.1	298.2	298.2/299.8	307.4/309.0
Br4-Br5	232.02(3)	233.3/237.7	239.8	233.2/235.1	232.8/234.8
Br1-Br6	290.71(3)	·		·	•
Br6-Br7	233.52(3)				
Br1-Br8	303.97(3)				
Br8-Br9	232.76(3)				
Br1-Br10	295.27(3)				
Br10-Br11	233.02(3)				
Br11-Br12	328.84(3)				
Br12-Br13	227.31 (5) ^[c]				
Br3-Br12	366.86(4)				
Br5-Br7	334.88(3)				

[a] For atom numbering, see Figure 2. [b] Square-pyramidal structures computed at DFT-D3 and ab initio level using the def2-TZVPP and *def-SV(P)* basis set. [c] Br13 is disordered; only the values for the major domain are discussed.

However, if we apply an ab initio method, such as SCS-MP2, which performed very well in previous investigations of polyhalides, including $[Br_{11}]^-$, we observe the D_{3h} and C_{4v} structures to be lower in energy (see the Supporting Information for reference [12]). This discrepancy is based on two facts: the lack of dispersion in the DFT treatment of polyhalogens, and the direct Hartree-Fock (HF) exchange in the DFT treatment (Figure 4). It is also seen that there is almost no energetic preference between the trigonal-bipyramidal and squarepyramidal structure, which is in accordance with other known fivecoordinate complexes. However,

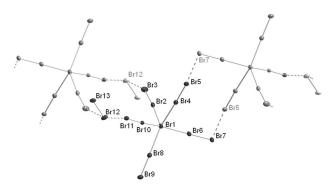


Figure 2. Chain-like structure of [PPN][Br $_{11}$ ·Br $_{2}$] in which the [Br $_{11}$] $^-$ units are connected by embedded Br $_{2}$ molecules. Ellipsoids are set at 50% probability and cations are omitted for clarity.

According to the procedure described by Addison et al., [19] the coordination of the five ligand Br_2 molecules around the central Br^- is almost square pyramidal in the present compound (Figure 1). The geometric parameter $\tau = (\beta - \alpha)/60^{\circ}$, where β and α are the largest angles in the coordination sphere, is determined to be 0.18 (for square pyramidal geometries $(C_{4\nu})$ $\tau = 0$; for trigonal pyramidal geometries (D_{3h}) $\tau = 1$). For more structural details, such as angles, see Table S7 in the Supporting Information.

Based on the previous gas-phase investigation of polyiodides, we compare our observed crystal structure of the $[Br_{11}]^-$ with the computed structure of $[I_{11}]^-$. It was found at the B3LYP level that the undecaiodide structure has C_s symmetry (Figure 3, **A**). This is in contrast to our obtained crystal structure, which shows an almost square-pyramidal arrangement of the Br_2 units around the central Br^- , which is also in agreement with the quantum-chemical results of Pichierri. Indeed if we optimize the most probable structures in C_s , D_{3h} , or $C_{4\nu}$ symmetry by DFT at the B3LYP level, we observe that the previously found structure of the $[I_{11}]^-$ in C_s symmetry is the lowest in energy for both undecahalides of iodine and bromine (Figure 3).

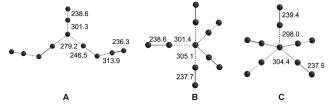


Figure 3. Computed $[Br_{11}]^-$ structures at the SCS-MP2/def2-SV(P) level. **A**: C_s symmetry, **B**: D_{3h} symmetry, and **C**: $C_{4\nu}$ symmetry.

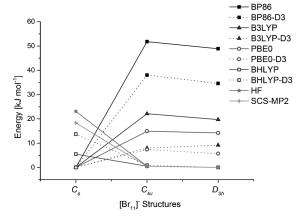


Figure 4. Energy diagram of different quantum-chemical methods applied to the three minima for the structures of $[Br_{11}]^-$ in C_s , $C_{4\rho}$ and D_{3h} symmetry.

the intermolecular exchange between these two isomers by Berry pseudo rotation is hindered in the solid. In this case the crystal packing effects force the molecular structure of the $[Br_{11}]^-$ anion to be nearly square-pyramidal (see above). This is in line with the systematic investigation of the influence of the cation size to the crystal structure of the polybromide anion.^[11]

As we have shown before, Raman spectroscopy is the method of choice for the investigation of such polyha-

lides.[12,16,22] Our experimental spectrum shows four Raman bands above 200 cm⁻¹, of which three bands agree well with our quantum-chemical calculations at SCS-MP2 level (Table 2 and Figure 6). The fifth band (v_1) is not found in

Table 2: Measured and computed Raman frequencies $[cm^{-1}]$ of $[Br_{11} \cdot Br_2]$ and molecular Br2.

Vibration	Exp.	SCS-MP2/ def2-TZVPP	SCS-MP2/ def-SV(P) Structure B	
		Structure C		
ν (Br _{2g})	325 ^[21]	334	310	
ν (Br _{2s})	308			
v_1 (Br ₂)	297			
ν_2	286	295	277	
ν_3	269	278	259	
ν_4	264	271	255	
$\nu_{\scriptscriptstyle 5}$		95	108	

our computed spectrum, it is assigned to the stretching mode of the embedded bromine molecule (Br12-Br13), which was not considered in our quantum-chemical description. The wavenumber of 297 cm⁻¹ is in fair agreement with experimentally observed solid bromine (308 cm⁻¹; Table 2). The vibrational modes v_{2-4} in the range 286–264 cm⁻¹ correspond to the stretching modes of the coordinated Br₂ units. Both modes v_3 and v_4 are close to each other, so that we only observe a shoulder for v_4 on v_3 in the experimental spectra. The v_5 band is the deformation mode of Br⁻···Br₂ (Figure 5).

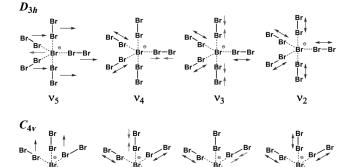


Figure 5. The most intense Raman modes of the $[Br_{11}]^-$ anion at D_{3h} and $C_{4\nu}$ symmetry.

 ν_3

These modes have been computed to be almost 170 cm⁻¹ below the Br₂ stretching modes of which the v_5 mode is by far the most intense. This v_5 band in the experimental spectrum is overlaid by the longitudinal and transversal vibrations of the solid, which usually appear below 150 cm⁻¹. In any case, our experimental and computed spectra show good agreement in the 300-200 cm⁻¹ region and support our experimental findings (Figure 6).

In conclusion, we reported herein the first structural evidence of a undecabromide [Br₁₁] stabilized by a bis(triphenylphosphine)iminium cation. It shows an almost square-

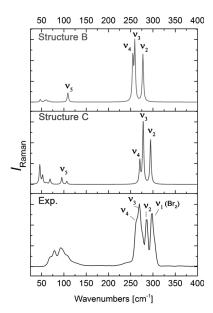


Figure 6. Experimental and computed (SCS-MP2) Raman spectra for [PPN][Br₁₁·Br₂] and [Br₁₁] $^-$, respectively.

pyramidal arrangment in the crystal structure. This outcome is supported by quantum-chemical calculations at ab-initio SCS-MP2 level, where the trigonal-pyramidal and also the squarepyramidal structure were shown to be the lowest in energy, indicating a very shallow potential energy surface. Furthermore, we have shown that the description of polyhalides by DFT methods is problematic, which is most likely due to the lack of dispersion and direct HF exchange in the DFT treatment.

This novel anion is the highest known ratio of elements to charge of any structurally known polyhalide compound and is therefore an ideal benchmark molecule for the investigation of weak interactions.

Experimental Section

[PPN]Br was synthesized as previously described. [23,24] The synthesis of [PPN]+[Br11·Br2] was carried out in pure elemental bromine by adding or condensing bromine (Merck) to [PPN]Br in a 1:6 ratio, forming a red-brown liquid. After maintaining the reaction mixture at room temperature for several days, red-brown clusters of crystals of [PPN][Br_{11} · Br_2] were obtained. At room temperature, the crystals are only stable in the reaction mixture.

The FT-Raman spectra were recorded on a Bruker Vertex 70 spectrometer equipped with a RAM II module using a Ge detector cooled with liquid nitrogen. Raman spectra were recorded (backscattering mode) at room temperature and also cooled with liquid nitrogen in flame-sealed glass capillaries (1064 nm, 10 mW power, resolution 4 cm⁻¹).

Crystal data for $[PPN]^+[Br_{11} Br_2]^-: C_{36}H_{30}Br_{13}N_1P_2, M_w = 1577.38,$ triclinic, space group $P\bar{1}$, a = 12.5230(2), b = 12.6557(2); c =16.1176(3) Å, $\alpha = 104.5740(10)$, $\beta = 106.0260(10)$, $\gamma = 96.4340(10)^{\circ}$, $V = 2330.48(7) \text{ Å}^3, Z = 2, \rho_{\text{calcd}} = 2.248 \text{ Mg m}^{-3}, F(000) = 1476, \lambda = 1476$ $0.71073 \text{ Å}, T = 100(2) \text{ K}, absorption coefficient} = 11.266 \text{ mm}^{-1},$ absorption correction: multi-scan, $T_{\min} = 0.5003$, $T_{\max} = 0.7469$. Data for the structure were collected on a Bruker SMART APEX2 CCD area-detector diffractometer with Mo K_{α} radiation. A single crystal has been taken directly from the reaction mixture coated at 0°C with



perfluoroether oil and mounted on a 0.1 mm Micromount. The structure was solved by direct methods in SHELXTL^[25] and OLEX2^[26] and refined by least squares on weighted F^2 values for all reflections. The final refinements converged at GooF = 1.024, R1 = 0.0312, and wR2 = 0.0621 for reflections ($I > 2\sigma(I)$), R1 = 0.0506, and wR2 = 0.0667 for all reflections. The hydrogen atoms were included in the refinement in calculated positions by a riding model. The graphical representations were prepared with Diamond. [27] CCDC 906333 ([PPN]+[Br₁₁·Br₂]-) contains the supplementary crystallographic data for this paper. These data can be obtained free of charge from The Cambridge Crystallographic Data Centre via www.ccdc.cam.ac.uk/data request/cif.

Quantum-chemical calculations were performed at various levels of density functional theory (DFT) and at ab initio HF, SCS-MP2 level. The gradient-corrected BP86^[28-32] functional, the hybrid functionals B3LYP,^[28-31,33,34] PBE0,^[28,29,35-37] and BHLYP^[28,29,31,33,38] with 25 % and 50 % HF exchange admixture were used. The def-SV(P) and def2-TZVPP basis sets were used for bromine and iodine. All calculations were carried out with the Turbomole V6.3^[39] program and the analytical gradient methods implemented therein. Structures have been fully optimized at DFT, HF, and SCS-MP2^[40] level. The Grimme dispersion correction D3 as implemented in Turbomole was used.^[41] Minima on the potential energy surface were characterized by harmonic vibrational frequency analyses, using numerical second derivatives based on energies and analytical gradients. We provide relative energies without zero-point vibrational corrections, as these do not alter the thermochemistry significantly.

Received: December 12, 2012 Revised: January 28, 2012 Published online: March 25, 2013

Keywords: bromine · halogens · polyhalogen anions · quantum-chemical calculations · Raman spectroscopy

- [1] S. M. Jørgensen, J. Prakt. Chem. 1870, 2, 347.
- [2] P. H. Svensson, L. Kloo, Chem. Rev. 2003, 103, 1649.
- [3] M. Groessl, Z. Fei, P. J. Dyson, S. A. Katsyuba, K. L. Vikse, J. S. McIndoe, *Inorg. Chem.* 2011, 50, 9728.
- [4] K. O. Strømme, Acta Chem. Scand. 1959, 13, 2089.
- [5] K. N. Robertson, P. K. Bakshi, T. S. Cameron, O. Knop, Z. Anorg. Allg. Chem. 1997, 623, 104.
- [6] C. W. Cunningham, G. R. Burns, V. McKee, *Inorg. Chim. Acta* 1990, 167, 135.
- [7] M. Wolff, J. Meyer, C. Feldmann, Angew. Chem. 2011, 123, 5073; Angew. Chem. Int. Ed. 2011, 50, 4970.
- [8] M. Wolff, A. Okrut, C. Feldmann, Inorg. Chem. 2011, 50, 11683.
- [9] H. Haller, S. Riedel, Nachr. Chem. 2012, 60, 865.

- [10] F. D. Chattaway, G. Hoyle, J. Chem. Soc. 1923, 123, 654.
- [11] X. Y. Chen, M. A. Rickard, J. W. Hull, C. Zheng, A. Leugers, P. Simoncic, *Inorg. Chem.* 2010, 49, 8684.
- [12] H. Haller, M. Ellwanger, A. Higelin, S. Riedel, Angew. Chem. 2011, 123, 11732; Angew. Chem. Int. Ed. 2011, 50, 11528.
- [13] V. Vitske, H. Herrmann, M. Enders, E. Kaifer, H.-J. Himmel, Chem. Eur. J. 2012, 18, 14108.
- [14] T. Chivers, J. F. Richardson, N. R. M. Smith, *Inorg. Chem.* 1985, 24, 2453.
- [15] J. Taraba, Z. Zak, Inorg. Chem. 2003, 42, 3591.
- [16] R. Brückner, H. Haller, M. Ellwanger, S. Riedel, *Chem. Eur. J.* 2012, 18, 5741.
- [17] S. Riedel, T. Köchner, X. F. Wang, L. Andrews, *Inorg. Chem.* 2010, 49, 7156.
- [18] B. Vonnegut, B. E. Warren, J. Am. Chem. Soc. 1936, 58, 2459.
- [19] A. W. Addison, T. N. Rao, J. Reedijk, J. van Rijn, G. C. Verschoor, J. Chem. Soc. 1984, 1349.
- [20] F. Pichierri, Chem. Phys. Lett. 2011, 515, 116.
- [21] K. P. Huber, G. Herzberg, *Molecular Spectra and Molecular Structure*, Vol. 4, Van Nostrand, New York, **1979**.
- [22] H. Haller, M. Ellwanger, A. Higelin, S. Riedel, Z. Anorg. Allg. Chem. 2012, 638, 553.
- [23] A. Martinsen, J. Songstad, Acta Chem. Scand. Ser. A 1977, 31, 645.
- [24] C. Knapp, R. Uzun, Acta Crystallogr. Sect. E 2010, 66, o3186.
- [25] G. M. Sheldrick, Acta Crystallogr. Sect. A 2008, 64, 112.
- [26] O. V. Dolomanov, L. J. Bourhis, R. J. Gildea, J. A. K. Howard, H. Puschmann, J. Appl. Crystallogr. 2009, 42, 339.
- [27] K. Brandenburg, *Diamond* 3.1 ed., Crystal Impact GbR, Bonn, 2009.
- [28] P. A. M. Dirac, Proc. R. Soc. London Ser. A 1929, 123, 714.
- [29] J. C. Slater, Phys. Rev. 1951, 81, 385.
- [30] S. H. Vosko, L. Wilk, M. Nusair, Can. J. Phys. 1980, 58, 1200.
- [31] A. D. Becke, Phys. Rev. A 1988, 38, 3098.
- [32] J. P. Perdew, Phys. Rev. B 1986, 33, 8822.
- [33] C. Lee, W. Yang, R. G. Parr, Phys. Rev. B 1988, 37, 785.
- [34] A. D. Becke, J. Chem. Phys. 1993, 98, 5648.
- [35] J. P. Perdew, Y. Wang, Phys. Rev. 1992, 45, 13244.
- [36] J. P. Perdew, K. Burke, M. Ernzerhof, Phys. Rev. Lett. 1996, 77, 3865.
- [37] J. P. Perdew, M. Ernzerhof, K. Burke, J. Chem. Phys. 1996, 105, 9982.
- [38] A. D. Becke, J. Chem. Phys. 1993, 98, 1372.
- [39] Turbomole 6.3 ed., a development of University of Karlsruhe and Forschungszentrum Karlsruhe GmbH, Karlsruhe, 2011, available from http://www.turbomole.com.
- [40] S. Grimme, J. Chem. Phys. 2003, 118, 9095.
- [41] S. Grimme, J. Antony, S. Ehrlich, H. Krieg, J. Chem. Phys. 2010, 132, 154104.